

Pr. Find the phase velocity and the magnitude of the attenuation constant of plane waves at a frequency 10^6 Hz in polyethylene, given (5)

$$\mu = \mu_0, \quad \epsilon_r = 2.3 \quad \text{and} \quad \sigma = 2.56 \times 10^{-4} \frac{\text{A}}{\text{m}}$$

Phase velocity

$$v = \frac{1}{\sqrt{\mu\epsilon}} = \frac{\sqrt{\epsilon_0/\epsilon}}{\sqrt{\mu_0\epsilon_0}} = c\sqrt{\epsilon_0/\epsilon}$$

$$= c\sqrt{\frac{1}{\epsilon_r}} = 3 \times 10^8 \times \sqrt{\frac{1}{2.3}}$$

$$v = 1.97 \times 10^8 \text{ m/s}$$

Since $\frac{\sigma}{\omega\epsilon} = \frac{2.56 \times 10^{-4}}{2\pi \times 10^6 \times 2.3 \times 8.85 \times 10^{-12}} = 2 \times 10^{-4}$

is much smaller than 1

$$\text{So } \beta = \omega\sqrt{\frac{\mu\epsilon}{2}} \left[\frac{\sigma^2}{2\epsilon^2\omega^2} \right]^{1/2}$$

$$= \frac{1}{2} \sigma \sqrt{\frac{\mu}{\epsilon}} = \frac{1}{2} \frac{\sigma}{\omega\epsilon} \cdot \omega\sqrt{\mu\epsilon}$$

$$= \frac{1}{2} \frac{\sigma}{\omega\epsilon} \cdot \frac{\omega}{v}$$

$$= \frac{1}{2} \times 2 \times 10^{-4} \times 2\pi \times 10^6 \times \frac{1}{1.97 \times 10^8}$$

$$\beta = 3.18 \times 10^{-2}$$

Reflection from the surface of a conductor (metal)

We extend the previous arguments to a boundary surface of a conducting medium.

We consider the simple case of normal incidence.

For incident, reflected and transmitted waves, (6)

$$E_I = E_{0I} \exp\{i(k_I r - \omega t)\}; H_I = \frac{k_I \times E_I}{\omega \mu_1} \quad (1)$$

$$E_R = E_{0R} \exp\{i(k_I r - \omega t)\}; H_R = \frac{k_I \times E_R}{\omega \mu_1} \quad (2)$$

$$E_T = E_{0T} \exp\{i(k_T r - \omega t)\}; H_T = \frac{k_T \times E_T}{\omega \mu_2} \quad (3)$$

Since the medium 2 is a conducting medium the propagation vector k_T is given by

$$k_T^2 = \epsilon_2 \mu_2 \omega^2 \left[1 + \frac{i\sigma}{\epsilon_2 \omega} \right]$$

The boundary conditions require that

$$E_{0I} - E_{0R} = E_{0T} \quad (4)$$

$$k_I (E_{0I} + E_{0R}) = k_T E_{0T} \quad (5)$$

Since k_T is complex, E_{0R} and E_{0T} cannot both be real hence one should expect phase shifts other than 0 or π in the reflected and transmitted waves.

Solving eq (4) and (5), we have

$$E_{0R} = \frac{k_T - k_I}{k_T + k_I} E_{0I} \quad (6)$$

$$E_{0T} = \frac{2k_I}{k_T + k_I} E_{0I}$$

substituting $k_I = \omega(\epsilon_1 \mu_1)^{1/2}$
and k_T

$$E_{OR} = \frac{\sqrt{\epsilon_2 \mu_2 \omega^2} \left(1 + \frac{i\sigma}{\epsilon_2 \omega}\right)^{1/2} - \omega(\epsilon_1 \mu_1)^{1/2}}{\sqrt{\epsilon_2 \mu_2 \omega^2} \left(1 + \frac{i\sigma}{\epsilon_2 \omega}\right)^{1/2} + \omega(\epsilon_1 \mu_1)^{1/2}} E_{OI} \quad (7)$$

and

$$E_{OT} = \frac{2\omega(\epsilon_1 \mu_1)^{1/2}}{\sqrt{\epsilon_2 \mu_2 \omega^2} \left(1 + \frac{i\sigma}{\epsilon_2 \omega}\right)^{1/2} + \omega(\epsilon_1 \mu_1)^{1/2}} E_{OI} \quad (8)$$

Let us now examine the case of a perfect conductor for which $\sigma = \infty$, then

$$E_{OR} = E_{OI} \quad \text{and} \quad E_{OT} = 0$$

hence, the ~~reflection~~ reflection is complete.

If the conductor is not a perfect conductor but a very good conductor $\frac{\sigma}{\epsilon_2 \omega} \gg 1$.

The approximation gives

$$k_T = \alpha + i\beta = (1+i) \sqrt{\frac{\omega \sigma \mu_2}{2}} = \frac{1+i}{\delta}$$

we have made use of the definition of skin depth δ .

Substituting this in eqⁿ (6)

$$E_{OR} = \frac{\frac{1+i}{\delta} - \omega(\epsilon_1 \mu_1)^{1/2}}{\frac{1+i}{\delta} + \omega(\epsilon_1 \mu_1)^{1/2}} E_{OI}$$

$$E_{OR} = \frac{\left\{ \frac{1}{\delta} - \omega(\epsilon_1 \mu_1)^{1/2} \right\} + \frac{i}{\delta}}{\left\{ \frac{1}{\delta} + \omega(\epsilon_1 \mu_1)^{1/2} \right\} + \frac{i}{\delta}} E_{OI}$$

Therefore, the reflection coefficient R is given by

$$R = \frac{|E_{0r}|^2}{|E_{0i}|^2} = \frac{1 - \omega(\epsilon_1 \mu_1)^{1/2} \delta^2 + 1}{1 + \omega(\epsilon_1 \mu_1)^{1/2} \delta^2 + 1}$$

Because $\frac{\sigma}{\epsilon_2 \omega} \gg 1$, $\omega(\epsilon_1 \mu_1)^{1/2} \delta \ll 1$, hence

$$R \approx 1 - 2\omega(\epsilon_1 \mu_1)^{1/2} \delta \quad \text{--- (9)}$$

$$= 1 - 2 \sqrt{\frac{2\omega \epsilon_1 \mu_1}{\sigma \mu_2}} \quad \text{--- (10)}$$

If the magnetic permeabilities are assumed to be equal

$$R = 1 - 2 \sqrt{\frac{2\omega \epsilon_1}{\sigma}} \quad \text{--- (11)}$$

The measure of energy transmitted into the conducting medium is obtained by calculating the transmission coefficient T

$$T = 1 - R = 2 \sqrt{\frac{2\omega \epsilon_1}{\sigma}} \quad \text{--- (12)}$$

Magnitude of the fraction of energy transmitted in a good conductor, we calculate T

for copper

$$\sigma = 6 \times 10^7 \text{ (ohm}^{-1}\text{m)}, \quad \nu = 10^{10} \text{ sec}^{-1}$$

$$T = 2 \sqrt{\frac{2 \times 2\pi \times 10^{10} \times 8.85 \times 10^{-12}}{6 \times 10^7}}$$

$$\approx 3 \times 10^{-4}$$

Which is extremely small for direct measurement to be made